

Lecture 9:

Unitary Braided fusion categories

$$\mathcal{C} = (L, N_c^{ab}, \underbrace{F_a^{abc}, R_c^{ab}})$$

$\leftarrow |L| =$
rank \mathcal{C}

Unitary matrices

$\mathcal{C} \sim$ Topological order

(Wen, Kas, etc.)

1-form symmetries

↳ gauging!

↳ Capelli

Analogous to F-symbols, R-symbols
are not gauge invariant

Given two UBFC

$\mathcal{C} \rightsquigarrow \bar{\mathcal{C}}$ Verifying whether \mathcal{C} and $\bar{\mathcal{C}}$

are gauge equivalent is a hard

task $\sim O(|L|)$.

Gauge invariant quantities: Modular data

Spoiler: Modular data does not exhaust the list of gauge invariant quantities in UBFC.

↳ Smallest example for this has rank 49.

Mod. data $\mathcal{E}_{49} = \text{Mod. data } \overline{\mathcal{E}}_{49}$

but $\mathcal{E}_{49} \not\rightarrow \overline{\mathcal{E}}_{49}$

Not gauge
equivalent.

Modular data: T and S matrices.

Consider UBFC $\mathcal{E} = (L, N_c^{ab}, F_a^{abc}, R_c^{ab})$

• for each $a \in L$ we define $d_a \in \mathbb{R}_{>0}$: quantum dimension (computed from N_c^{ab} , measures "non-Abelianness" of a).

• for each $a \in L$ we define $\Theta_a \in U(1)$: topological spin of a .

(computed from R -symbols, measures self-braiding)

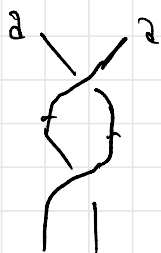
Defined $\bigcirc_a \uparrow = \Theta_a \big|_a \uparrow$ or twist

no non-Abelian anyon $d_a \neq 1$ \leftrightarrow unitary matrices capture operations in this space

$\Rightarrow \Theta_a = \frac{1}{d_a} \sum_c d_c \text{Tr}[\mathcal{R}_{c, a}^{aa}]$ ArXiv:1506.05805

\uparrow Ho-weak.

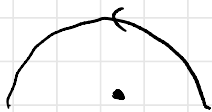
\uparrow trace over $p, v = 1, \dots, N_c^{aa}$



a : non-Abelian anyons.

$(\mathcal{R}_{c, a}^{aa})^2 \sim$ braiding matrix

$$\Psi(x_2, x_1) = U \Psi(x_1, x_2)$$



$$|\Theta_a| = 1$$

$$\left\{ \begin{array}{l} \Theta_a = 1 \sim \text{boson} \Rightarrow s_a = 0 \\ \Theta_a = -1 \sim \text{fermion} \Rightarrow s_a = 1/2 \\ \Theta_a = e^{i\pi/2} \sim \text{semion} \Rightarrow s_a = 1/4 \end{array} \right.$$

$$\Theta_a = e^{2\pi i s_a}$$

$$s_a = \arg \frac{\Theta_a}{2\pi} \text{ mod } 1$$

\hookrightarrow topological spin.

for general sa no anyon
 ↳ "any" top. spin

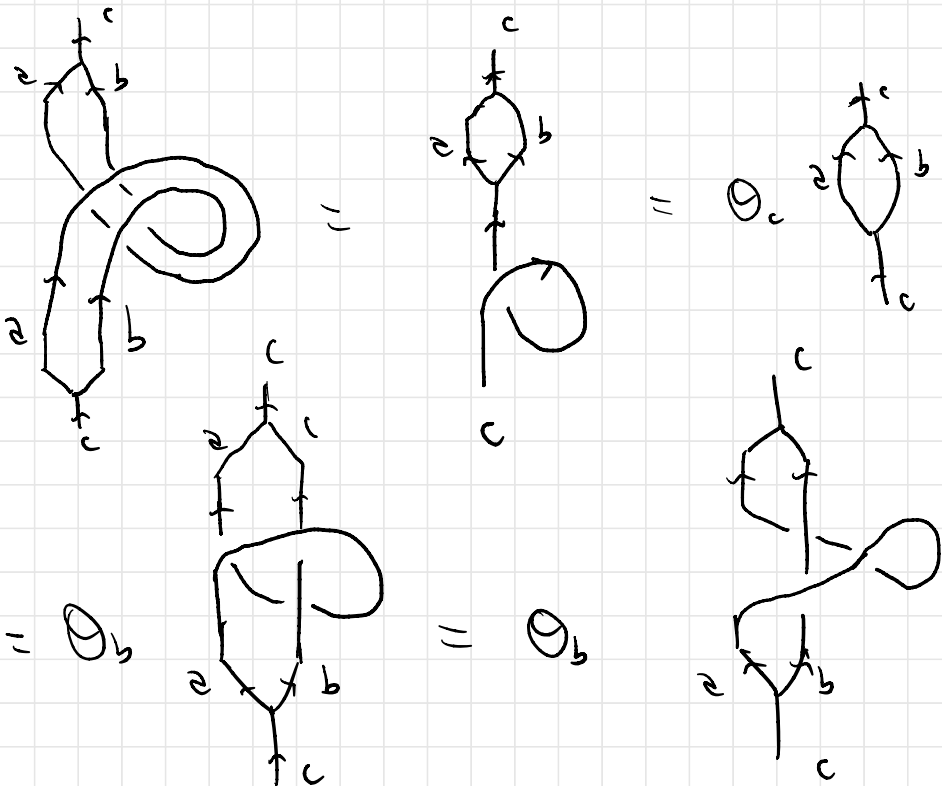
⇒ One can show that

$$R_{\pm}^{a a^*} = \Theta_a^* \chi_a$$

↳ Specified by $\text{sgn}(F_{i,j,k})$
 to be checked.

↳ ± 1 Frobenius-Schur indicator of $a \in \mathcal{L}$.

They obey Ribbon equation.



Similarly, we diagrammatically define $|L| \times |L|$ matrix

$$S_{ab} = \frac{1}{D} a \left(\bigcirc \bigcirc \right) b$$

$$D = \sqrt{\sum_a d_a^2}$$

One can show (Homework)

$$S_{ab} = \frac{1}{D} \sum_c N_c^{ab} \text{Tr} (R_c^{ab} R_c^{ba}) d_c$$

↑ matrix indices
 $\mu, \nu = 1, \dots, N_c^{ab}$

$$S_{ab} = \frac{1}{D} \sum_c N_c^{ab} \frac{\partial_c}{\partial_a \partial_b} d_c$$

These S and T matrices must obey

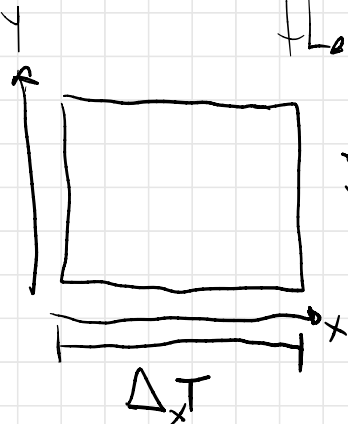
$$S^2 \equiv C \sim \text{charge-conjugation matrix.}$$

$$(ST)^3 \equiv e^{\frac{2\pi i c_-}{8}} C$$

S & T matrices projective representations of modular group $SL(2, \mathbb{Z})$

c_- : chiral central charge

- $c_- \neq 0 \Rightarrow$ chiral CFT at boundary.
- $c_- \neq 0 \Rightarrow$ very robust edge modes
 or itrary perturbations
 [Levin 1301.7355]
- Experimentally $c_- \neq 0$ implies in a thermal Hall conductance K_{xy}



$$J_y = K_{xy} \Delta_x T$$

$$K_{xy} = \frac{c_- \pi^2 k_B^2}{3h} T.$$

- C_- is only defined mod 8

$$C_- \rightarrow C_- + 8$$

$\hookrightarrow \exists$ topological phase E_8 with no non-trivial anyons in the bulk, but non-trivial boundary w/ $C_- = 8$

$\hookrightarrow E_8 \sim$ "Hidden topological phase"

$\Rightarrow E_8$ level WZW model

$\Rightarrow K$ -matrix CS Theory.

$$K = \begin{pmatrix} 2 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 1 & 2 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 2 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 2 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 2 & 1 & 0 & 1 \\ 0 & 0 & 0 & 0 & 1 & 2 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 2 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 2 \end{pmatrix} \quad \det K = 1$$

$$\text{sgn } K = 8$$

$\text{sgn } K = \# \text{ positive eigenvalues} - \# \text{ negative eigenvalues}$

$\hookrightarrow 16$ -fold way classification (invertible bosonic phases)

$\hookrightarrow e \in L = \{L\}$

\Rightarrow Stacking

Consider the set S of all topological orders.

|| We can equip this set with \boxtimes (Deligne product) which is a stacking operation

$$\boxed{e} \boxtimes \boxed{D} = \boxed{e \boxtimes D}$$

$$e = (L, N_c^{ab}, F_d^{abc}, R_c^{ab})$$

$$D = (L', N_c'^{ab}, F_d'^{abc}, R_c'^{ab})$$

$$\rightarrow \underbrace{e \boxtimes D}_{\text{"new topological order"}} = (L \times L', N_c^{ab} N_c'^{ab}, F_d^{abc} \otimes F_d'^{abc}, R_c^{ab} \otimes R_c'^{ab})$$

$$G \stackrel{\text{local unitary}}{\cong} e \boxtimes D.$$

$\rightarrow e_{\text{triv}} = (\{1\}, n_i = 1, F=1, R=1)$
behaves as unit element.

$\Rightarrow S$ with \boxtimes is Monoid. (group with elements not necessarily having inverses).

Eg: $C = \text{Toric}$ \mathbb{R}^4 , $D = \text{Fib.}$ \mathbb{R}^2

$\Rightarrow C \boxtimes D$
 \mathbb{R}^6 anyons

there exists a subset A of S whose elements all have inverses. A forms an Abelian group

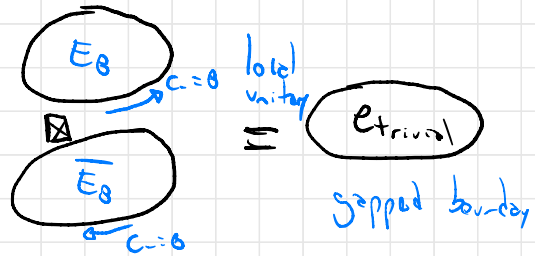
top. order $\in A \Rightarrow$ Invertible topological orders.

\hookrightarrow for all $e \in A$, there exists an inverse \bar{e} such that

$$e \boxtimes \bar{e} \equiv e_{\text{triv}}$$

Necessary condition for $e \in A \Rightarrow |e| = 1$.

E.g.: • E_0 state



• IQHE

• Topological superconductors,

⇒ Inevitable phases: non-trivial edge modes with no protecting symmetries; trivial bulk (no anyons).

Ex: $\mathcal{E} = \text{Fib}^{\delta^2}$, $\mathcal{D} = \text{Ising}^{\delta^3}$
 $\mathcal{E} \boxtimes \mathcal{D}$ has 6 anyons $a \boxtimes b$
 \mathcal{E}_{Fib} $\mathcal{E}_{\text{Ising}}$

E.g.:

$\mathcal{E} \boxtimes \mathcal{D} \supset \tau \boxtimes \sigma$ Consider Ising:

$$\begin{aligned}
 (\tau \boxtimes \sigma) \otimes (\tau \boxtimes \sigma) &= (\tau \otimes \tau) \boxtimes (\sigma \otimes \sigma) \\
 &= (1 + \tau) \boxtimes (1 + \psi)
 \end{aligned}$$

$$= | \boxtimes | + | \boxtimes \psi + \tau \boxtimes | + \tau \boxtimes \psi$$